

Computational Fluids Mechanics

Part 1: Numerical methods

Part 2: Turbulence models

Part 3: Practice of numerical simulation

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Computational Fluid Mechanics

Part1 : Numerical Methods

Objective :

Numerical methods for fluids mechanics and their accuracy

1. Introduction
2. Numerical methods for fluid mechanics (1) ~ Basic equations
3. Numerical methods for fluid mechanics (2) ~Discretizing schemes
4. Numerical methods for fluid mechanics (3) ~Coupling algorism
5. Numerical methods for fluid mechanics (4) ~Additional problems
6. Reliability of numerical simulation

Summary

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Computational Fluid Mechanics

Part1 : Numerical Methods (2a)

1. Introduction

2. Numerical methods for fluid mechanics (1) ~ governing equations

- Governing equations of fluid flow
- Typical solutions of fluid flow
- Additional models for complex flow phenomena

3. Numerical methods for fluid mechanics (2) ~Discretizing schemes

4. Numerical methods for fluid mechanics (3) ~Coupling algorism

5. Numerical methods for fluid mechanics (4) ~Additional problems

6. Reliability of numerical simulation

Governing Eqs. of Fluid mechanics

Continuity Eq.
(mass conservation)

$$\frac{D\rho}{Dt} = -\rho I$$

Momentum eq.

$$\rho \frac{Du_i}{Dt} = -\frac{\partial p}{\partial x_i} + \frac{\mu}{3} \frac{\partial I}{\partial x_i} + \mu \nabla^2 u_i + \rho f_i$$

Energy Eq.

$$\rho c_v \frac{DT}{Dt} = -pI + \lambda \nabla^2 T + \Phi + \dot{q}$$

Eq. of state

$$p = \rho RT$$

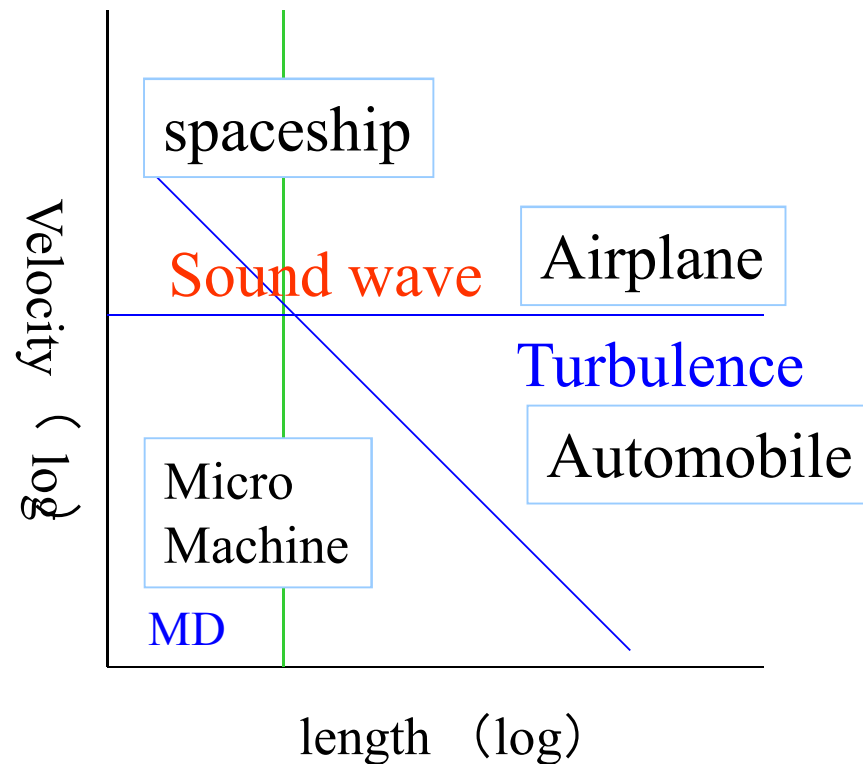
$$; \frac{D}{Dt} \equiv \frac{\partial}{\partial t} + u_j \frac{\partial}{\partial x_j} \quad ; \quad \nabla^2 \equiv \frac{\partial^2}{\partial x_k \partial x_k}$$

$$; I = \frac{\partial u_k}{\partial x_k} \quad ; \quad \Phi = \frac{\partial u_i}{\partial x_j} \left(\frac{\mu}{3} I \delta_{ij} + \mu \frac{\partial u_i}{\partial x_j} \right)$$

Governing Eqs. of Fluid Flow

Non-dimensional Parameters

- **Reynolds no.**
$$\frac{(\text{length}) \times (\text{velocity})}{(\text{viscosity})}$$
- **Mach no.**
$$(\text{velocity}) / (\text{Sound speed})$$
- **Knussen no.**
$$(\text{free pass}) / (\text{length})$$



Non-dimensional Parameters

- in the governing eq.-

$$\left[\frac{L}{t_0 U} \right] \frac{\partial \rho^*}{\partial t^*} + u_j^* \frac{\partial \rho^*}{\partial x_j^*} = -\rho^* \frac{\partial u_j^*}{\partial x_j^*} \quad p^* = \rho^* T^* \quad \frac{p_0}{\rho_0 T_0} = R$$

$$\left[\frac{L}{t_0 U} \right] \rho^* \frac{\partial u_i^*}{\partial t^*} + \rho^* u_j^* \frac{\partial u_i^*}{\partial x_j^*} = - \left[\frac{p_0}{\rho_0 U^2} \right] \frac{\partial p^*}{\partial x_i^*} + \left[\frac{\mu}{\rho_0 U L} \right] \left(\frac{1}{3} \frac{\partial I^*}{\partial x_i^*} + \frac{\partial^2 u_i^*}{\partial x_j^{*2}} \right) + \left[\frac{F}{U^2 L} \right] \rho^* f_i^*$$

$$\left[\frac{L}{t_0 U} \right] \rho^* \frac{\partial T^*}{\partial t^*} + \rho^* u_j^* \frac{\partial T^*}{\partial x_j^*} = - \left[\frac{p_0}{\rho_0 c_v T_0} \right] p^* I^* + \left[\frac{\lambda}{\rho_0 c_v U L} \right] \frac{\partial^2 T^*}{\partial x_j^{*2}} + \left[\frac{\mu}{\rho_0 U L} \right] \left[\frac{U^2}{c_v T_0} \right] \Phi^* + \left[\frac{L Q}{\rho_0 c_v U} \right] \dot{q}^*$$

$$t = t^* t_0, \quad x_j = x_j^* L, \quad u_i = u_i^* U, \quad p = p^* p_0, \quad T = T^* T_0, \quad \rho = \rho^* \rho_0$$

$$f = F f^*, \quad \dot{q} = Q \dot{q}^*$$

Non-dimensional Parameters

- in the governing eq.-

Parameter of fluid flow

$$\left[\frac{L}{t_0 U} \right] = St \quad (\text{Strouhal no.})$$

$$\left[\frac{p_0}{\rho_0 c_v T_0} \right] = \frac{\rho_0 R T_0}{\rho_0 c_v T_0} = \gamma - 1$$

(γ : specific heat ratio)

$$\left[\frac{p_0}{\rho_0 U^2} \right] = \frac{a_0^2}{\gamma U^2} = \frac{1}{\gamma M^2}$$

(M: Mach No)

$$\left[\frac{\lambda}{\rho_0 c_v U L} \right] = \text{Pe} = \text{Re} \cdot \text{Pr}$$

(Pe: Peclet No., Pr: Prandtl No.)

$$\left[\frac{\mu}{\rho_0 U L} \right] = \frac{1}{\text{Re}} \quad (\text{Reynolds No.})$$

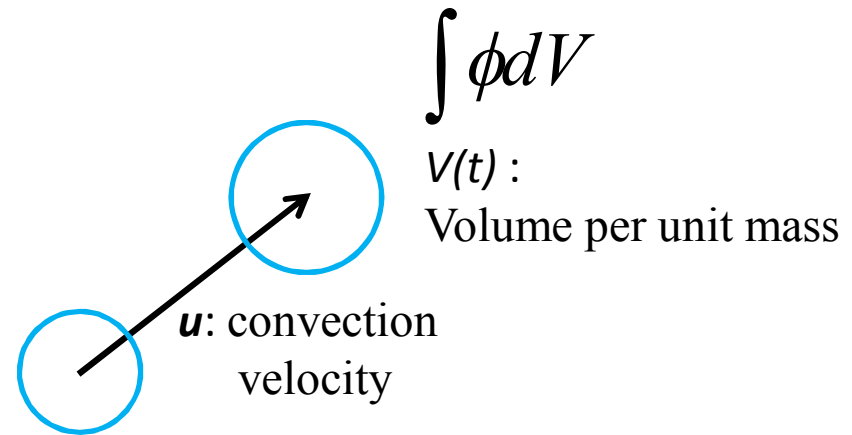
$$\left[\frac{U^2}{c_v T_0} \right] = (\gamma - 1) M^2$$

Lagrangean vs. Eularian formulations

Lagrangean

$$\frac{D\phi}{Dt} \equiv \left(\frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla \phi \right)$$

on unit mass [$*/\text{kg}$]



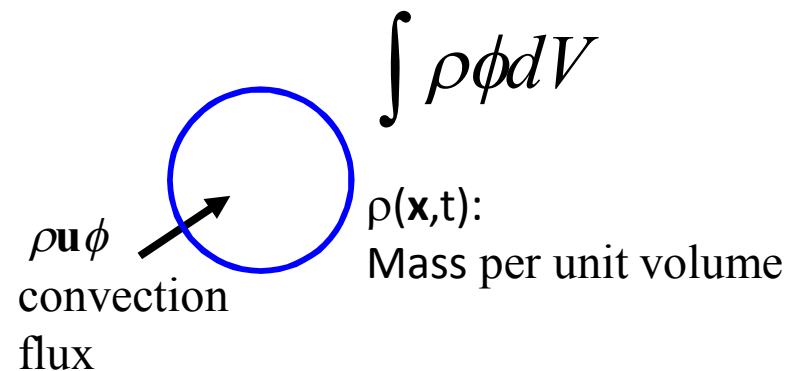
Eularian

$$\frac{\partial \rho \phi}{\partial t} + \nabla \cdot \rho \mathbf{u} \phi$$

$$\equiv \rho \frac{D\phi}{Dt} + \phi \left(\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} \right)$$

$= 0$ (Mass conservation)

on unit volume [$*/\text{m}^3$]



Conservation law

Mass flux vector

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{u})$$

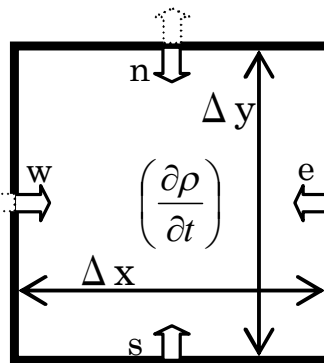
$$\frac{\partial \rho}{\partial t} = -\left(\frac{\partial \rho u}{\partial x} + \frac{\partial \rho v}{\partial y}\right)$$

$$\int \frac{\partial \rho}{\partial t} dV = -\int (\rho \mathbf{u}) \cdot \mathbf{n} dS$$

$$\left(\frac{\partial \rho}{\partial t}\right) \Delta x \Delta y \cong (\rho u_w - \rho u_e) \Delta y + (\rho v_s - \rho v_n) \Delta x$$

$$\begin{aligned} -(\rho \mathbf{u} \cdot \mathbf{n})_n &= -(\rho u, \rho v) \cdot (0, 1) \\ &= -\rho v_n \end{aligned}$$

$$\begin{aligned} -(\rho \mathbf{u} \cdot \mathbf{n})_w &= \\ -(\rho u, \rho v) \cdot (-1, 0) &= \rho u_w \end{aligned}$$



$$\begin{aligned} -(\rho \mathbf{u} \cdot \mathbf{n})_e &= -(\rho u, \rho v) \cdot (1, 0) \\ &= -\rho u_e \end{aligned}$$

$$\begin{aligned} -(\rho \mathbf{u} \cdot \mathbf{n})_s &= -(\rho u, \rho v) \cdot (0, -1) \\ &= \rho v_s \end{aligned}$$

Conservation law

Flux vector vs. Stress tensor

$$\frac{\partial \rho \mathbf{u}^t}{\partial t} = -\nabla \cdot (\rho \mathbf{u} \mathbf{u}^t) - \nabla \cdot \mathbf{T} - \nabla \cdot \mathbf{P} + \mathbf{f}^t \quad \Rightarrow \quad \frac{\partial \rho u}{\partial t} = -\nabla \cdot (\rho \mathbf{u} u) - \nabla \cdot \boldsymbol{\tau}_x - \nabla \cdot \mathbf{p}_x + f_x$$

$$\rho \mathbf{u}^t = \begin{bmatrix} \rho u & \rho v \end{bmatrix}, \quad \mathbf{f}^t = \begin{bmatrix} f_x & f_x \end{bmatrix} \quad \rho \mathbf{u} \mathbf{u}^t = \begin{bmatrix} \rho u u & \rho u v \\ \rho v u & \rho v v \end{bmatrix},$$

$$\mathbf{P} = \begin{bmatrix} \mathbf{p}_x & \mathbf{p}_y \end{bmatrix} = \begin{bmatrix} p & 0 \\ 0 & p \end{bmatrix} (= p \mathbf{I} = p \delta_{ij}) \quad \therefore \nabla \cdot \mathbf{P} = (\nabla p)^t = \begin{bmatrix} \frac{\partial p}{\partial x} & \frac{\partial p}{\partial y} \end{bmatrix}$$

$$\mathbf{T} = \begin{bmatrix} \boldsymbol{\tau}_x & \boldsymbol{\tau}_y \end{bmatrix} = \begin{bmatrix} -2\mu \frac{\partial u}{\partial x} & -\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \\ -\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) & -2\mu \frac{\partial v}{\partial y} \end{bmatrix} (= 2\mu \mathbf{S})$$

Conservation law

Momentum flux vector

$$\frac{\partial \rho u}{\partial t} = -\nabla \cdot (\mathbf{J}_{conv.} + \mathbf{J}_{diff.} + \mathbf{J}_{pres.}) + f_x$$

$$\mathbf{J}_{conv.} = \rho \mathbf{u} u = \begin{bmatrix} \rho u u \\ \rho v u \end{bmatrix}, \quad \mathbf{J}_{diff.} = \begin{bmatrix} -2\mu(\partial u / \partial x) \\ -\mu(\partial u / \partial y + \partial v / \partial x) \end{bmatrix}, \quad \mathbf{J}_{pres.} = \begin{bmatrix} p \\ 0 \end{bmatrix}$$

$$\int \frac{\partial \rho u}{\partial t} dV = -\int (\mathbf{J}_{conv.} + \mathbf{J}_{diff.} + \mathbf{J}_{pres.}) \cdot \mathbf{n} dS + \int f_x dV$$

$$-(\mathbf{J}_{conv.} \cdot \mathbf{n})_w = J_{conv.} = (\rho u u)_w$$

$$-(\mathbf{J}_{diff.} \cdot \mathbf{n})_w = J_{diff.} = -2\mu \frac{\partial u}{\partial x}_w$$

$$-(\mathbf{J}_{pres.} \cdot \mathbf{n})_w = J_{pres.} = p_w$$

$$-(\mathbf{J}_{conv.} \cdot \mathbf{n})_e = -J_{conv.} = -(\rho u u)_e$$

$$-(\mathbf{J}_{diff.} \cdot \mathbf{n})_e = -J_{diff.} = 2\mu \frac{\partial u}{\partial x}_e$$

$$-(\mathbf{J}_{pres.} \cdot \mathbf{n})_e = -J_{pres.} = -p_e$$

Exercise 2a

Calculate each component of momentum flux through the vertical surface (n & s in fig.) of considering volume.

Derive the conservation law for y-direction component of momentum (ρv)

Computational Fluid Mechanics

Part1 : Numerical Methods (2b)

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 - **Typical solutions of fluid flow**
 - Additional models for complex flow phenomena
3. Numerical methods for fluid mechanics (2) ~Discretizing schemes
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6. Reliability of numerical simulation

Typical solutions of fluid flow

convection

$$\frac{\partial \rho}{\partial t} = -u_j \frac{\partial \rho}{\partial x_j}$$

$$\rho \frac{\partial u_i}{\partial t} = -\rho u_j \frac{\partial u_i}{\partial x_j}$$

$$\rho c_v \frac{\partial T}{\partial t} = -\rho c_v u_j \frac{\partial T}{\partial x_j}$$

interaction (volume source)

$$-\rho \frac{\partial u_k}{\partial x_k}$$

$$-\frac{\partial p}{\partial x_i} + F_i$$

$$-p \frac{\partial u_k}{\partial x_k} + Q + \Phi$$

$$p = \rho RT$$

diffusion (surface flux)

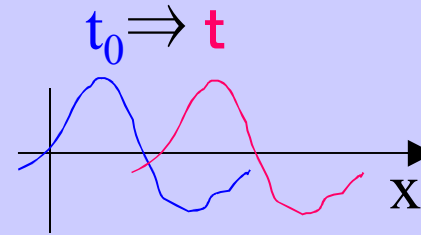
$$+\frac{\partial}{\partial x_j} \left(\frac{\mu}{3} \frac{\partial u_k}{\partial x_k} \delta_{ij} + 2\mu S_{ij} \right)$$

$$+\frac{\partial}{\partial x_j} \left(k \frac{\partial T}{\partial x_j} \right)$$

Solution of convection

- Model of scalar convection

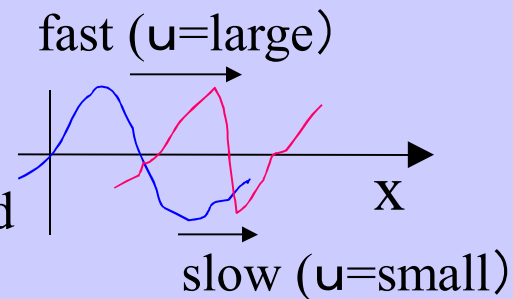
(linear eq.)
$$\frac{\partial \phi}{\partial t} = -u \frac{\partial \phi}{\partial x} \Rightarrow \phi(x, t) = \phi_0(x - ut)$$



- Model of momentum convection

(Burger's Eq.)
$$\frac{\partial u}{\partial t} = -u \frac{\partial u}{\partial x}$$

- shape deformed
- high frequency produced



- ➔
- keep linear solution
 - adaptive to steep variation

Solution of diffusion and external force

Diffusion of scalar (Linear)

time dependent heat transfer problem

$$\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2} \quad \Rightarrow \quad \phi = \phi_n \exp\left(inx - \frac{t}{\Gamma}\right)$$

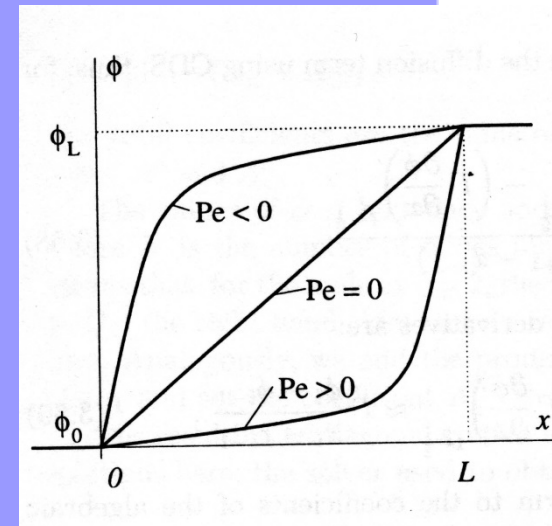
steady convective – diffusion problem

$$u \frac{\partial \phi}{\partial x} = \Gamma \frac{\partial^2 \phi}{\partial x^2} \quad \Rightarrow \quad \left(u\phi - \Gamma \frac{\partial \phi}{\partial x} \right) = \text{const}$$

Pelet No.: $Pe = uL/\Gamma$

Source term (Linear)

$$\frac{\partial \phi}{\partial t} = -S\phi \quad \Rightarrow \quad \phi = \phi_0 \exp(-St)$$



Solution of Interaction step

Compressible (sound wave)

$$\frac{\partial \rho}{\partial t} = -\rho \frac{\partial u}{\partial x}$$

$$\rho \frac{\partial u}{\partial t} = -\frac{\partial p}{\partial x}$$

$$\frac{p}{\rho^\kappa} = \text{const}$$

(S=const.)



$$\frac{\partial \rho}{\partial t} = -\rho \frac{\partial u}{\partial x}$$

$$\rho \frac{\partial u}{\partial t} = -a^2 \frac{\partial \rho}{\partial x}$$

$$a^2 \equiv \left(\frac{dp}{d\rho} \right)_s = \sqrt{\kappa \frac{p}{\rho}}$$

Mach No. : $Ma = U / a$
 a: sound speed
 U: flow velocity

$$\frac{\partial^2 u}{\partial t^2} = a^2 \frac{\partial^2 u}{\partial x^2} \quad \text{Sound wave}$$

$$\Rightarrow u = f(x - at) + g(x + at)$$

Solution of Interaction step

Incompressible flow

$$\nabla \cdot \mathbf{u} = \frac{\partial u_i}{\partial x_i} = 0$$

$$\frac{\partial u_i}{\partial t} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + F_i$$

$$(F_i = -u_j \frac{\partial u_i}{\partial x_j} + \dots)$$

$$\rho = \text{const}$$



if $\text{Ma} \rightarrow 0$
 a (sound speed) $\rightarrow \infty$

$$\frac{1}{\rho} \nabla \cdot \nabla p = -\frac{\partial}{\partial t} (\nabla \cdot \mathbf{u}) + \nabla \cdot \mathbf{F}$$

Pressure poisson eq.

Solution of Interaction step

Compressible (1D nozzle)

$$\left. \frac{d\rho}{\rho} + \frac{du}{u} + \frac{dA}{A} = 0 \right\}$$

$$u du + \frac{dp}{\rho} = 0$$

$$\left. \frac{dp}{p} - \kappa \frac{d\rho}{\rho} = 0 \right\}$$

$$\left. \frac{dp}{p} = \frac{d\rho}{\rho} + \frac{dT}{T} \right\}$$

$$\frac{d\rho}{\rho} = -\frac{M^2}{M^2 - 1} \frac{dA}{A}$$

$$\frac{du}{u} = +\frac{1}{M^2 - 1} \frac{dA}{A}$$

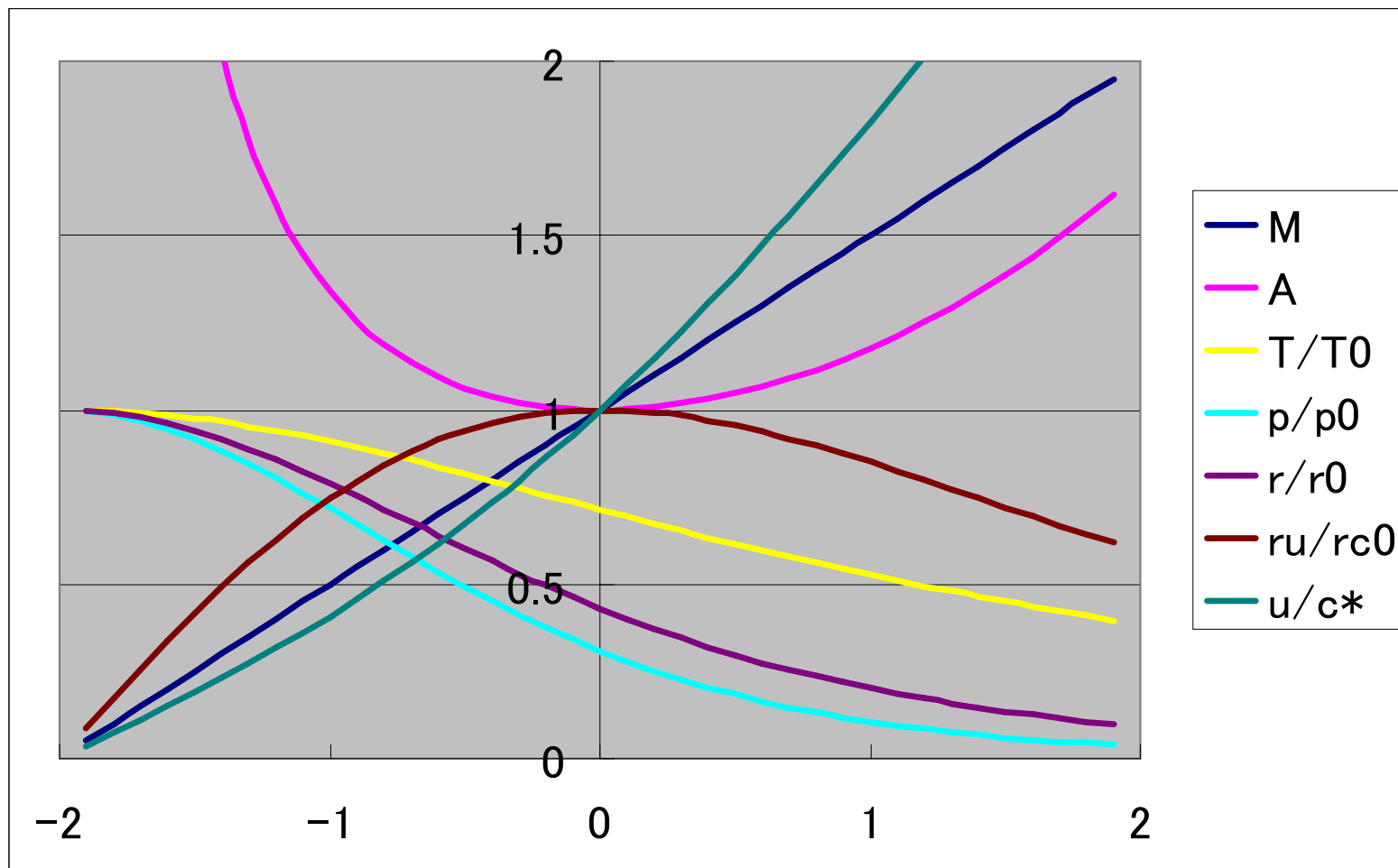
$$\frac{dp}{p} = -\frac{\kappa M^2}{M^2 - 1} \frac{dA}{A}$$

$$\left. \frac{dT}{T} = -\frac{(\kappa - 1)M^2}{M^2 - 1} \frac{dA}{A} \right\}$$

$$\frac{da}{a} = -\frac{(\kappa - 1)M^2}{2(M^2 - 1)} \frac{dA}{A}$$

$$\frac{dM}{M} = +\frac{2 + (\kappa - 1)M^2}{M^2 - 1} \frac{dA}{A}$$

Solution of 1D compressible flow



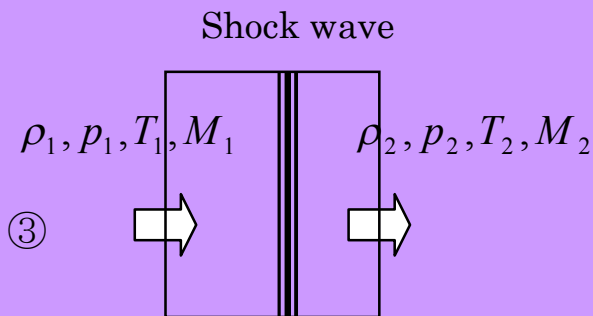
Solution of Interaction step

Compressible (1D shock wave)

$$\text{Mass} \quad : \quad \rho_1 u_1 = \rho_2 u_2 \quad \textcircled{1}$$

$$\text{Momentum} \quad : \quad p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2 \quad \textcircled{2}$$

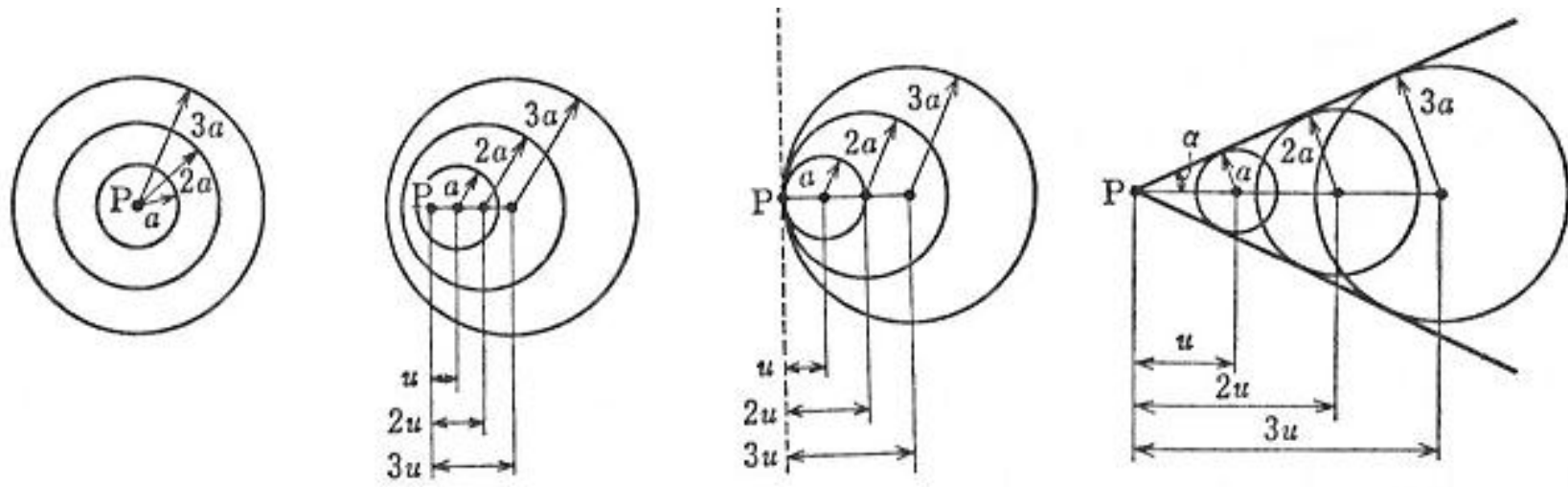
$$\text{Energy} \quad : \quad \frac{\kappa}{\kappa-1} \frac{p_1}{\rho_1} + \frac{1}{2} u_1^2 = \frac{\kappa}{\kappa-1} \frac{p_2}{\rho_2} + \frac{1}{2} u_2^2 \quad \textcircled{3}$$



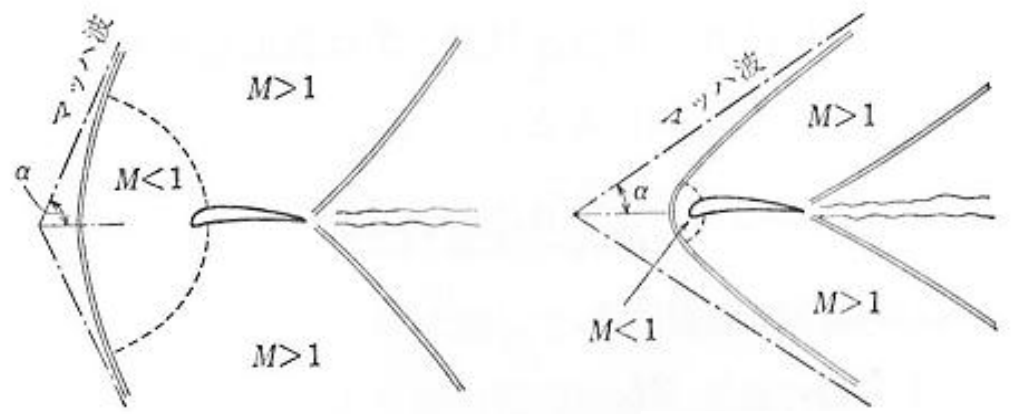
$$\therefore \frac{\rho_2}{\rho_1} = \frac{\left(\frac{\kappa+1}{\kappa-1} \frac{p_2}{p_1} + 1 \right)}{\left(\frac{p_2}{p_1} + \frac{\kappa+1}{\kappa-1} \right)} = \frac{u_1}{u_2} > 1 \quad \text{Rankine-Hugoniot eq.}$$

$$\therefore \frac{T_2}{T_1} = \frac{p_2}{p_1} \frac{\rho_1}{\rho_2} = \frac{p_2}{p_1} \frac{\left(\frac{\kappa+1}{\kappa-1} \frac{p_2}{p_1} + 1 \right)}{\left(\frac{p_2}{p_1} + \frac{\kappa+1}{\kappa-1} \right)} = \frac{\left(\frac{p_2}{p_1} + \frac{\kappa+1}{\kappa-1} \right)}{\left(\frac{p_1}{p_2} + \frac{\kappa+1}{\kappa-1} \right)} > 1$$

Profiles of shock wave



(a) $u=0, M=0$ (b) $u < a, 0 < M < 1$ (c) $u = a, M = 1$ (d) $u > a, M > 1$



(d) M_∞ medium (e) M_∞ large

Exercise 2b

Solve the following equations of $\phi(x)$ and draw the profiles for different parameter values.

$$u \frac{\partial \phi}{\partial x} = \Gamma \frac{\partial^2 \phi}{\partial x^2} \quad u, \Gamma \text{ are constant parameter}$$

$$0 \leq x \leq L \quad \phi = 0 \quad \text{at } x = 0$$

$$\phi = 1 \quad \text{at } x = L$$

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Complexity of system

Strong coupling

	convection	interaction	diffusion	external-force	
$\frac{\partial \rho}{\partial t}$	$= -u_j \frac{\partial \rho}{\partial x_j}$	$- \rho \frac{\partial u_k}{\partial x_k}$	$+ M$		mass
$\rho \frac{\partial u_i}{\partial t}$	$= -\rho u_j \frac{\partial u_i}{\partial x_j}$	$- \frac{\partial p}{\partial x_i}$	$+ \frac{\partial}{\partial x_j} \left(\frac{\mu}{3} I \delta_{ij} + 2\mu S_{ij} \right) + F_i$		momentum
$\rho c_v \frac{\partial T}{\partial t}$	$= -\rho c_v u_j \frac{\partial T}{\partial x_j}$	$- p \frac{\partial u_k}{\partial x_k}$	$+ \frac{\partial}{\partial x_j} \left(k \frac{\partial T}{\partial x_j} \right) + \Phi + Q$		energy
		$p = \rho RT$			

1) hyperbolic (compressible)

2) harmonic (sound wave)

3) parabolic (incompressible)

Ma > 1

Ma < 1

**Combined solution
is needed**

Complexity of system

Weak coupling

convection interaction diffusion production

$\frac{\partial \rho}{\partial t}$	$= -u_j \frac{\partial \rho}{\partial x_j}$	$- \rho \frac{\partial u_k}{\partial x_k}$	$+ M$	mass
$\rho \frac{\partial u_i}{\partial t}$	$= -\rho u_j \frac{\partial u_j}{\partial x_j}$	$- \frac{\partial p}{\partial x_i}$	$+ \frac{\partial}{\partial x_j} \left(\frac{\mu}{3} I \delta_{ij} + 2\mu S_{ij} \right) + F_i$	momentum
$\rho c_v \frac{\partial T}{\partial t}$	$= -\rho c_v u_j \frac{\partial T}{\partial x_j}$	$- p \frac{\partial u_k}{\partial x_k}$	$+ \frac{\partial}{\partial x_j} \left(k \frac{\partial T}{\partial x_j} \right) + \Phi + Q$	energy

$p = \rho RT$

**Alternative solution
Is available**

Additional eqs.

- Turbulence model
- NonNewtonian fluid

- Buoyancy force
- Multi-phase model
- Chemical reaction

Model of Combustion Flow

– Basic eqs. of reactive flow–

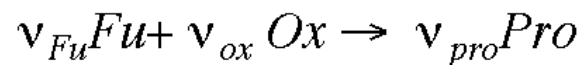
$$\begin{aligned}
 \text{Mass:} \quad & \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i} (\rho u_j) = 0 \\
 \text{Momentum:} \quad & \frac{\partial \rho u_i}{\partial t} + \frac{\partial}{\partial x_j} (\rho u_i u_j) = -\frac{\partial p}{\partial x_i} + \frac{\partial}{\partial x_j} (2\mu S_{ij}) + \frac{\partial}{\partial x_j} \left\{ \left(\mu_V - \frac{2}{3}\mu \right) \frac{\partial u_k}{\partial x_k} \delta_{ij} \right\} \\
 \text{Energy:} \quad & \rho C_p \frac{\partial T}{\partial t} + \rho C_p u_j \frac{\partial T}{\partial x_j} = \frac{\partial}{\partial x_j} \left[\frac{\mu}{Pr} \frac{\partial T}{\partial x_j} \right] + \frac{\partial p}{\partial t} \\
 \text{Eq. of gas state:} \quad & p = \rho R_o T \sum_{\alpha=1}^N Y_\alpha / M_\alpha + \left[\sum_{\alpha=1}^N C_{p_\alpha} \frac{\mu}{Sc_\alpha} \frac{\partial Y_\alpha}{\partial x_j} \right] \frac{\partial T}{\partial x_j} - \sum_{\alpha=1}^N h_{o_\alpha} \omega_\alpha
 \end{aligned}$$

+

$$\text{Chemical species} \quad \frac{\partial \rho Y_\alpha}{\partial t} + \frac{\partial}{\partial x_j} (\rho Y_\alpha u_j) = \frac{\partial}{\partial x_j} \left[\frac{\mu}{Sc_\alpha} \frac{\partial Y_\alpha}{\partial x_j} \right] + \omega_\alpha$$

Arrhenius' law

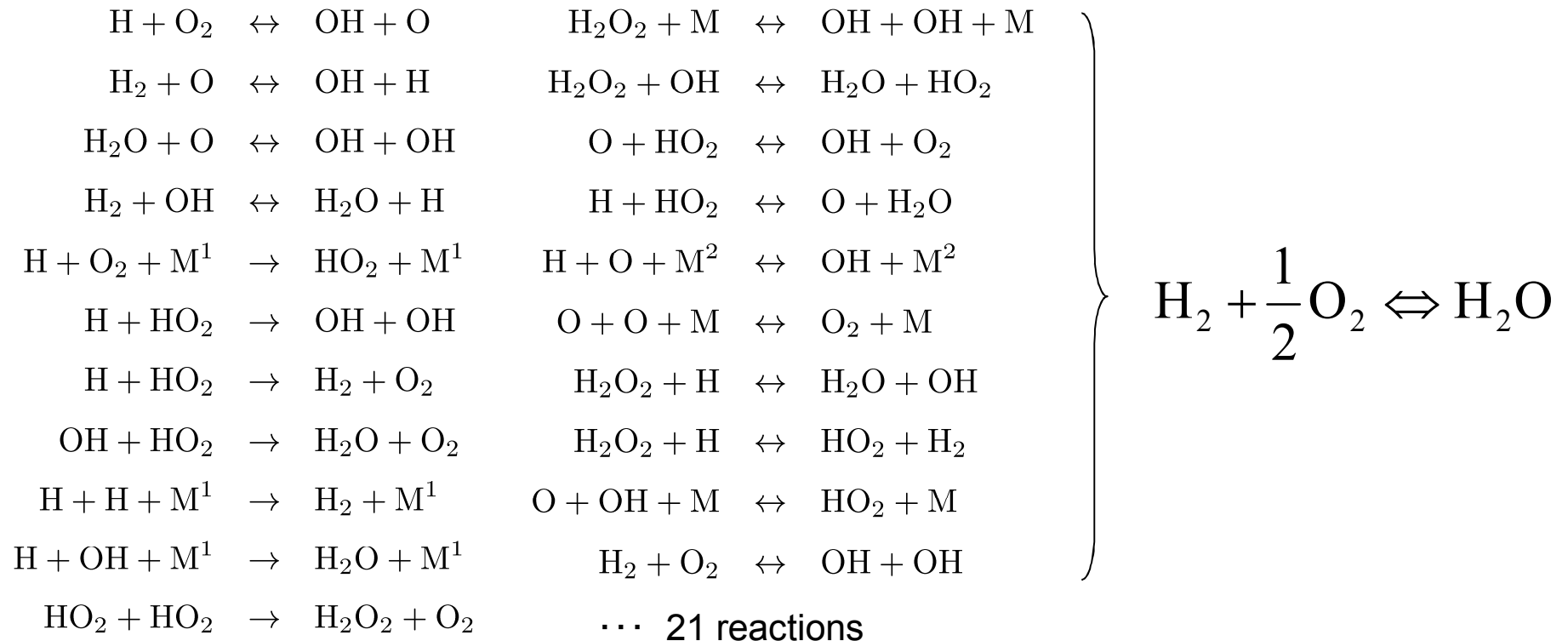
Chemical reaction



$$\omega_{fu} = -B \exp\left(-\frac{E}{RT}\right) \left(\frac{\rho Y_{fu}}{M_{fu}}\right)^{\nu_{fu}} \left(\frac{\rho Y_{ox}}{M_{ox}}\right)^{\nu_{ox}}$$

Detail reactions of H₂-O₂ flame

Spices	N ₂	O ₂	H ₂	H ₂ O	OH
Lewis No.	1	1.11	0.3	1.12	0.73
Spices	O	H	H ₂ O ₂	HO ₂	-
Lewis No.	0.7	0.18	1.12	1.1	-



Model of Combustion Flow

– non-dimensional parameters of mass & heat transfer–

- Nusselt no. $Nu \equiv \frac{h\lambda}{L}$ h: heat transfer coefficient
 λ : Thermal conductivity

- Prandtl no. $Pr \equiv \frac{\nu}{\alpha}$ ν : kinetic viscosity

- Schmidt no. $Sc \equiv \frac{\nu}{D}$ D: mass diffusion conductivity

- Lewis no. $Le \equiv \frac{\alpha}{D}$ α : temp. diffusion conductivity

Model of Combustion Flow

– non-dimensional parameters of reactive flow–

- (primal) Damköhler no. $D \equiv \frac{\tau_r}{\tau_c}$: time scale of flow
 τ_c : time scale of reaction

- Karlovits no. $K \equiv \frac{\eta}{U} \frac{dU}{dy}$ Thickness of heating zone: $\eta = \frac{\lambda}{\rho c_p S} = \frac{\alpha}{S}$
 Combustion speed: S

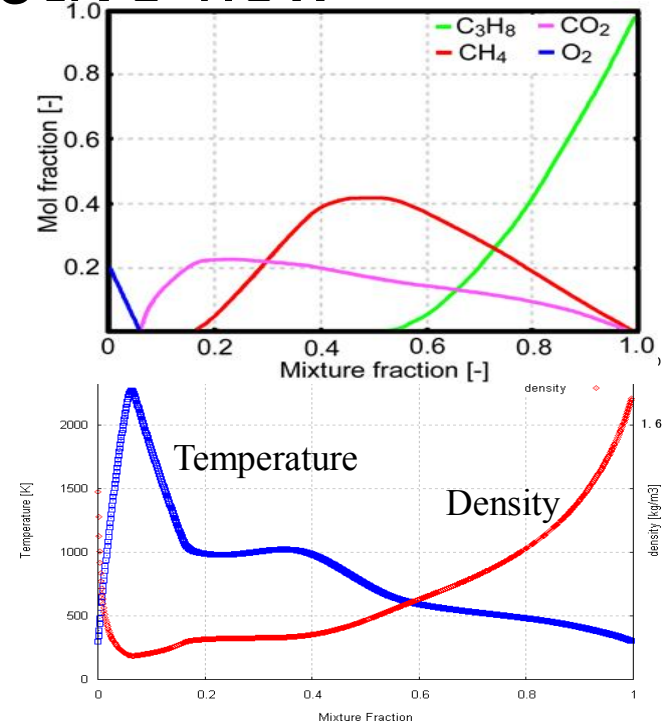
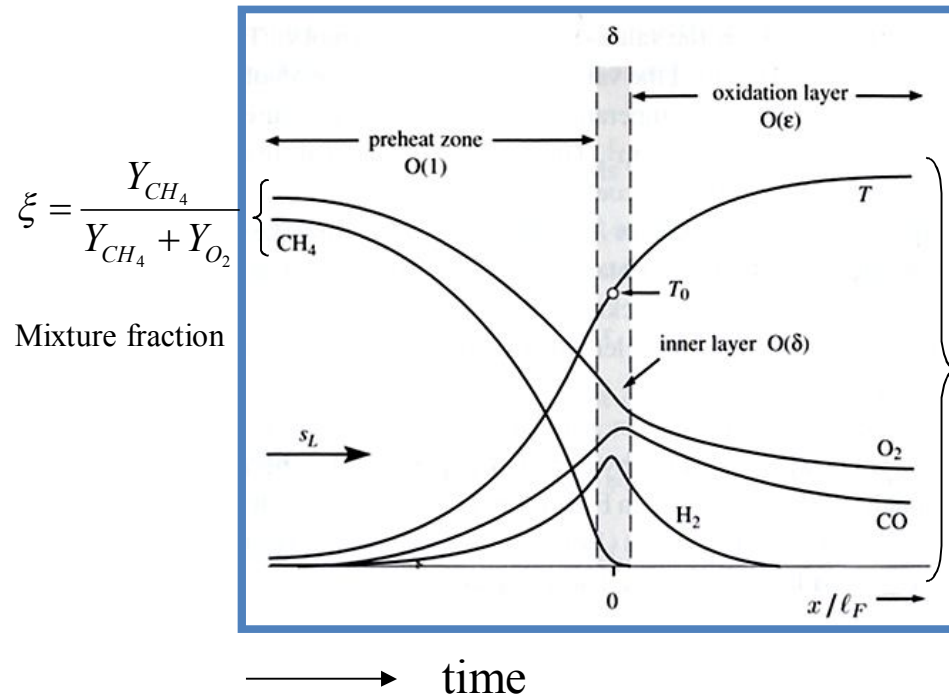
- turbulent Damköhler no.

$$Da \equiv \frac{\tau_t}{\tau_c} = \frac{l}{u' \delta}$$

l : turbulent length scale
 u' : velocity fluctuation, δ : flame thickness

Model of Combustion

– mechanics of reactive flow –



Enthalpy: $H = E + PV \Rightarrow dH = (TdS - PdV) + d(PV) = TdS + VdP$

$$\left(\frac{\partial H}{\partial S}\right)_p = T > 0$$

Free energy: $G = H - TS \Rightarrow dG = (TdS + VdP) - d(TS) = -SdT + VdP$

$$\left(\frac{\partial G}{\partial T}\right)_p = -S < 0$$

Enthalpy components:

$$dH = \sum_{\alpha} h_{\alpha}^0 dn_{\alpha} + c_p dT$$

$$\left(\frac{\partial H}{\partial T}\right)_p = c_p$$

Chemical potential + thermal energy

Mass transport

$$\frac{\partial \rho_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{V}_\alpha) = \dot{\omega}_\alpha$$



$$Y_\alpha = \rho_\alpha / \rho$$

$$\frac{\partial \rho Y_\alpha}{\partial t} + \nabla \cdot (\rho Y_\alpha \mathbf{V}_\alpha) = \dot{\omega}_\alpha$$



$$\mathbf{v}_{d\alpha} = \mathbf{V}_\alpha - \mathbf{V} \quad ; \quad \mathbf{V} = \frac{\sum \rho \mathbf{V}_\alpha}{\rho}$$

$$\rho \frac{DY_\alpha}{Dt} \equiv \rho \frac{\partial Y_\alpha}{\partial t} + \underbrace{\rho \mathbf{V} \cdot \nabla Y_\alpha}_{\text{convection}} = \underbrace{-\nabla \cdot (\rho Y_\alpha \mathbf{v}_{d\alpha})}_{\text{diffusion}} + \dot{\omega}_\alpha$$



$$\rho Y_\alpha \mathbf{v}_{d\alpha} \approx -\rho D_\alpha \nabla Y_\alpha \quad \text{Fick's law}$$

$$\rho \frac{DY_\alpha}{Dt} = \frac{\partial \rho Y_\alpha}{\partial t} + \nabla \cdot (\rho Y_\alpha \mathbf{V}) = \nabla \cdot (\rho D_\alpha \nabla Y_\alpha) + \dot{\omega}_\alpha \quad D_\alpha = \frac{\mu}{Sc_\alpha} \quad Le_\alpha = \frac{Pr_\alpha}{Sc_\alpha}$$

$$\rho = \sum \rho_\alpha, \quad \sum \dot{\omega}_\alpha = 0, \quad \sum Y_\alpha \mathbf{v}_{d\alpha} = 0 \quad \frac{D\rho}{Dt} \equiv \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Mass transport

Distribution of mole fraction

$$\nabla X_\alpha = \sum_{\beta} \frac{X_\alpha Y_\beta}{D_{\alpha\beta}} (\mathbf{v}_\beta - \mathbf{v}_\alpha) + (Y_\alpha - X_\alpha) \frac{\nabla p}{p}$$

$$X_\alpha = \frac{Y_\alpha / W_\alpha}{\sum_{\beta} Y_\beta / W_\beta}$$

$$+ \frac{\rho}{p} \sum_{\beta} Y_\alpha Y_\beta (\mathbf{f}_\alpha - \mathbf{f}_\beta) + \sum_{\beta} \frac{X_\alpha Y_\beta}{\rho D_{\alpha\beta}} \left(\frac{D_{T\beta}}{Y_\beta} - \frac{D_{T\alpha}}{Y_\alpha} \right) \frac{\nabla T}{T}$$

In case of two species

$$Y_1 \mathbf{v}_1 = -D_{12} \left\{ \nabla Y_1 + \frac{Y_1 Y_2}{X_1 X_2} (Y_1 - X_1) \frac{\nabla p}{p} + \frac{(Y_1 Y_2)^2}{X_1 X_2} \frac{\rho}{p} (\mathbf{f}_1 - \mathbf{f}_2) + \frac{D_{T,1}}{\rho D_{T,2}} \frac{\nabla T}{T} \right\}$$

Fick's law

$$Y_\alpha Y_\beta (\mathbf{v}_\beta - \mathbf{v}_\alpha) = Y_\alpha Y_\beta (\mathbf{V}_{d\beta} - \mathbf{V}_{d\alpha}) = Y_\alpha \underbrace{(Y_\beta \mathbf{V}_{d\beta})}_{= -Y_\alpha \mathbf{V}_{d\alpha}} - Y_\beta (Y_\alpha \mathbf{V}_{d\alpha}) = \underbrace{(Y_\alpha + Y_\beta)}_{= 1} (-Y_\alpha \mathbf{V}_{d\alpha})$$

Energy transport

Theory.

$$\rho dh = dp + \rho T ds$$

Enthalpy eq.

$$\rho \frac{Dh}{Dt} = \frac{Dp}{Dt} + \nabla \cdot \lambda \nabla h - \sum \nabla \cdot \rho D \sum (1 - Le_\alpha) \nabla Y_\alpha + \Phi + Q_R$$

Temperature eq.

$$\rho C_p \frac{DT}{Dt} = \frac{Dp}{Dt} + \nabla \cdot \lambda \nabla T + \sum c_{p,\alpha} \mathbf{j}_\alpha \cdot \nabla T + \Phi + Q_R + \dot{Q}_{reac}$$

$$dh = \sum Y_\alpha h_\alpha = \sum (h_\alpha^\circ dY_\alpha + c_{p,\alpha} Y_\alpha dT), \quad h_\alpha = h_\alpha^\circ + \int_{T_0}^T c_{p,\alpha} dT$$

Reactive heat

$$\sum h_\alpha^\circ \dot{\omega} = -\dot{Q}_{reac.} \quad \text{: Reactive heat}$$

$$C_p = \sum Y_\alpha C_{p,\alpha} \quad \text{Specific heat}$$

$$\alpha = \lambda / \rho C_p \quad \text{Thermal conductivity}$$

$$\text{Mass flux (Fick's law)} \quad \mathbf{j}_\alpha = -\rho D_\alpha \nabla Y_\alpha \quad Le_\alpha = \lambda / (\rho C_p D_\alpha) = \alpha / D_\alpha$$

$$\text{Low Ma No. approx.} \quad \nabla p \approx 0 \quad \therefore \frac{Dp}{Dt} \equiv \frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p \approx \frac{dp}{dt}$$



Solution of combustion flow

	convection	interaction	diffusion	external-force	
$\frac{\partial \rho}{\partial t}$	$= -u_j \frac{\partial \rho}{\partial x_j}$	$- \rho \frac{\partial u_k}{\partial x_k}$	$+ M$		mass
$\rho \frac{\partial u_i}{\partial t}$	$= -\rho u_j \frac{\partial u_i}{\partial x_j}$	$- \frac{\partial p}{\partial x_i}$	$+ \frac{\partial}{\partial x_j} \left(\frac{\mu}{3} I \delta_{ij} + 2\mu S_{ij} \right) + F_i$		momentum
$\rho c_v \frac{\partial T}{\partial t}$	$= -\rho c_v u_j \frac{\partial T}{\partial x_j}$	$- p \frac{\partial u_k}{\partial x_k}$	$+ \frac{\partial}{\partial x_j} \left(k \frac{\partial T}{\partial x_j} \right) + \Phi + Q$		energy
		$p = \rho RT$		↑ Heat source	
Additional eqs.		↑	Weak relation by low Ma approx.		
$\rho \frac{\partial Y_\alpha}{\partial t}$	$= -\rho u_j \frac{\partial Y_\alpha}{\partial x_j}$		$+ \frac{\partial}{\partial x_j} \left(D \frac{\partial Y_\alpha}{\partial x_j} \right) + \dot{\omega}_\alpha$		chemical species

Exercise 2c

Consider a system of equations for complex flow phenomena and analyze its coupling process to fluid dynamics;

Ex. Multi-phase flow

Buoyancy flow

Magneto-hydrodynamics

Flow in porous media

Plasma flow etc.